

# Matrix Decompositions and Quantum Circuit Design

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# Motivation

**Classical Problem:** Design AND-OR-NOT circuit for  $\varphi : (\mathbb{F}_2)^n \rightarrow \mathbb{F}_2$ , with  $\mathbb{F}_2 = \{0, 1\}$

**One Answer:** (see e.g. *Feynman on Computation*, section 2.4) Wire an AND circuit for each bit string on which  $\varphi = 1$ ; connect circuit blocks by OR's

- Restatement:
  - Produce a **decomposition** of the function  $\varphi$
  - Produce circuit blocks accordingly

**Matrix decompositions:** decompose unitary matrices,  
e.g. **quantum computations**

## Motivation, Cont.

However, the approach described here is so simple and general that it does not need an expert in logic to design it! Moreover, it is also a standard type of layout that can easily be laid out in silicon. (ibid.)

### Remarks:

- Analog for quantum computers?
- Simple & general?

# Outline

- I. Two Qubit Circuits (CD)
- II. Optimal Relative Phase Circuits
- III. Half CNOT per Entry (CSD)
- IV. Differential Topology & Lower Bounds

# The Magic Basis of Two-Qubit State Space

$$\begin{cases} |m0\rangle &= (|00\rangle + |11\rangle)/\sqrt{2} \\ |m1\rangle &= (|01\rangle - |10\rangle)/\sqrt{2} \\ |m2\rangle &= (i|00\rangle - i|11\rangle)/\sqrt{2} \\ |m3\rangle &= (i|01\rangle + i|10\rangle)/\sqrt{2} \end{cases}$$

**Remark:** Bell states up to global phase; global phases needed for theorem

**Theorem** (Lewenstein, Kraus, Horodecki, Cirac 2001)

Consider a  $4 \times 4$  unitary  $u$ , global-phase chosen for  $\det(u) = 1$

- Compute matrix elements in the magic basis
- (All matrix elements are real)  $\iff (u = a \otimes b)$

# Two-Qubit Canonical Decomposition

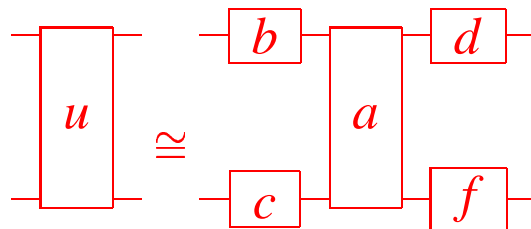
**Two-Qubit Canonical Decomposition:** Any  $u$  a four by four unitary admits a matrix decomposition of the following form:

$$u = (d \otimes f)a(b \otimes c)$$

for  $b \otimes c, d \otimes f$  are tensors of one-qubit computations,  $a = \sum_{j=0}^3 e^{i\theta_j} |mj\rangle \langle mj|$

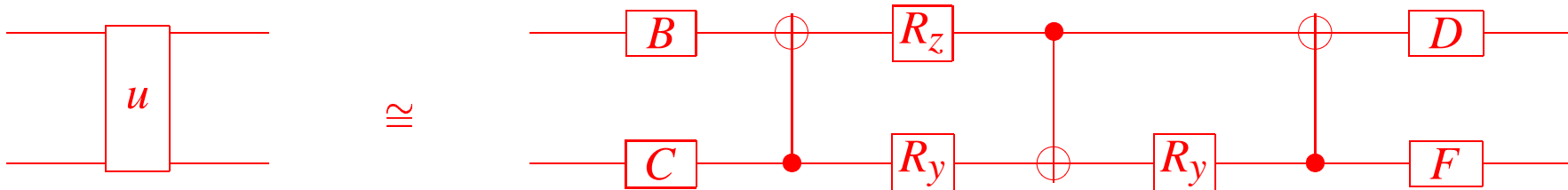
Note that  $a$  applies relative phases to the magic or Bell basis.

Circuit diagram: For any  $u$  a two-qubit computation, we have:



# Application: Three CNOT Universal Two-Qubit Circuit

- **Many groups: 3 CNOT circuit for  $4 \times 4$  unitary:**  
(F.Vatan, C.P.Williams), (G.Vidal, C.Dawson), (V.Shende, I.Markov, B-)
  - Implement  $a$  somehow, commute SWAP through circuit to cancel
  - Earlier B-,Markov: 4 CNOT circuit w/o SWAP, CD & naïve  $a$



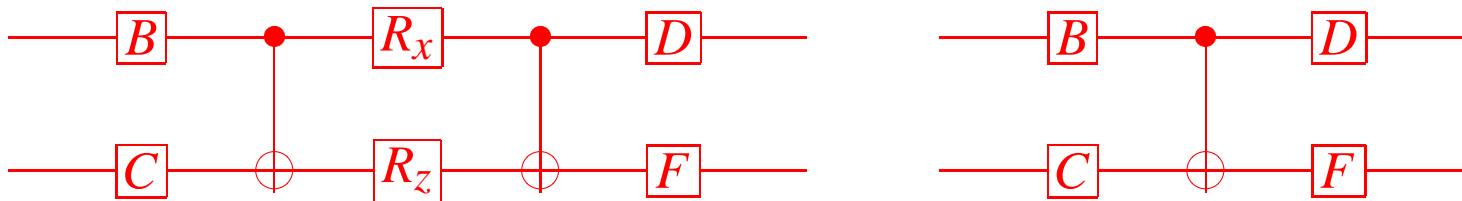
## Two-Qubit CNOT-Optimal Circuits

Theorem:(Shende,B-,Markov) Suppose  $u$  is a  $4 \times 4$  unitary **normalized so  $\det(u) = 1$** . Label  $\gamma(v) = (-i\sigma^y)^{\otimes 2}v(-i\sigma^y)^{\otimes 2}v^T$ . Then any  $v$  admits a circuit holding elements of  $SU(2)^{\otimes 2}$  and **3 CNOT's**, up to global phase. Moreover, for  $p(\lambda) = \det[\lambda I_4 - \gamma(v)]$  the characteristic poly of  $\gamma(v)$ :

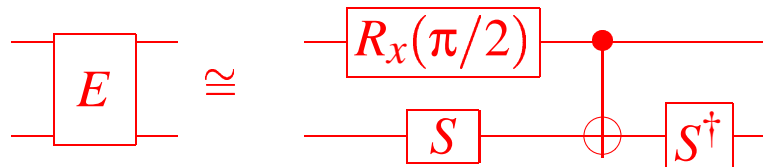
- ( $v$  admits a circuit with 2 CNOT's)  $\iff (p(\lambda)$  has real coefficients)
- ( $v$  admits a circuit with 1 CNOT)  $\iff (p(\lambda) = (\lambda + i)^2(\lambda - i)^2)$
- ( $v \in SU(2) \otimes SU(2)$ )  $\iff ( \gamma(v) = \pm I_4 )$



# Optimal Structured Two-qubit Circuits



- **Quantum circuit identities:** All 1,2 CNOT diagrams reduce to these
- Computing parameters: useful to use operator  $E$ ,  $E|j\rangle = |mj\rangle$



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# Relative Phase Group

- **Easiest conceivable  $n$ -qubit circuit question:** How to build circuits for

$$A(2^n) = \left\{ \sum_{j=0}^{2^n-1} e^{i\theta_j} |j\rangle\langle j| ; \theta_j \in \mathbb{R} \right\}?$$

- $A(2^N)$  commutative  $\implies$  vector group
  - $\log A(2^n) \rightarrow \mathfrak{a}(2^n)$  carries matrix multiplication to vector sum
  - Strategy: build decompositions from **vector space decompositions**
  - Subspaces encoded by **characters**, i.e. continuous group maps  $\chi : A(2^n) \rightarrow U(1)$

## Characters Detecting Tensors

- $\ker \log \chi$  is a subspace of  $\mathfrak{a}(2^n)$
- Subspaces  $\bigcap_j \ker \log \chi_j$  exponentiate to **closed** subgroups

**Example:**  $a = \sum_{j=0}^{2^n-1} z_j |j\rangle \langle j| \in A(2^n)$  has  $a = \tilde{a} \otimes R_z(\alpha)$  if and only if

$$z_0/z_1 = z_2/z_3 = \cdots = z_{2^n-2}/z_{2^n-1}$$

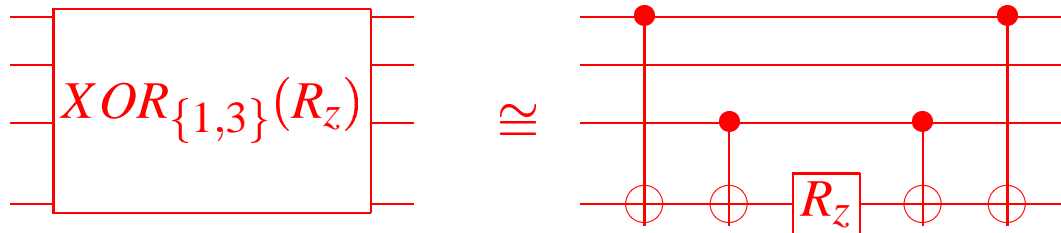
So  $a$  factors on the bottom line if and only if  $a \in \bigcap_{j=0}^{2^{n-1}-1} \ker \chi_j$   
for  $\chi_j(a) = z_{2j}z_{2j+2}/(z_{2j+1}z_{2j+3})$ .

## Circuits for $A(2^n)$

### Outline of Synthesis for $A(2^n)$ :

- Produce circuit blocks capable of setting all  $\chi_j = 1$
- After  $a = \tilde{a} \otimes R_z$ , induct to  $\tilde{a}$  on top  $n - 1$  lines

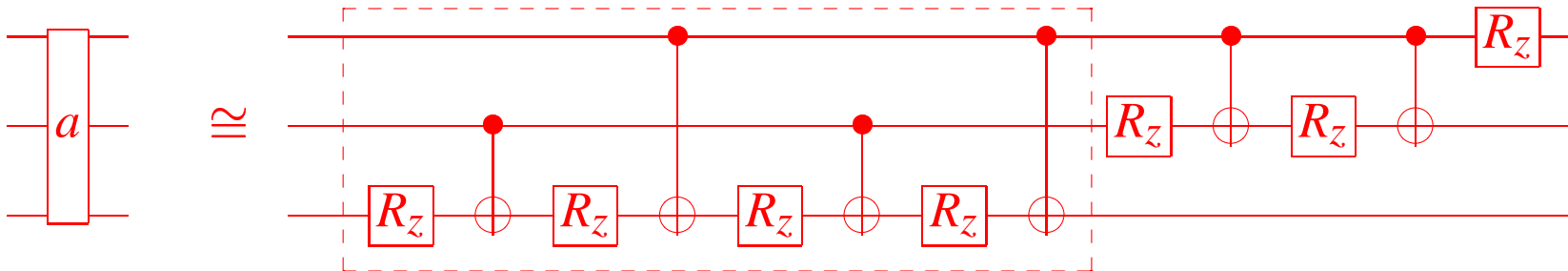
**Remark:**  $2^{n-1} - 1$  characters to zero  $\implies 2^{n-1} - 1$  blocks, i.e. one for each nonempty subset of the top  $n - 1$  lines



## Circuits for $A(2^n)$ , Cont.

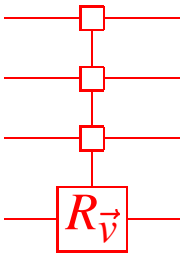
## Tricks in Implementing Outline:

- If  $\#[(S_1 \cup S_2) - (S_1 \cap S_2)] = 1$ , then all but one CNOT in center of  $XOR_{S_1}(R_z) XOR_{S_2}(R_z)$  cancel
- Take subsets in Gray code, most CNOTs cancel
- Final count:  $2^n - 2$  CNOTs



# Uniformly Controlled Rotations (M.Möttönen, J.Vartiainen)

Let  $\vec{v}$  be any axis on Bloch sphere. Uniformly-controlled rotation requires  $2^{n-1}$  CNOTs:

$$\bigwedge_k^{\text{uni}} [R_{\vec{v}}] = \begin{pmatrix} R_{\vec{v}}(\theta_0) & \mathbf{0}_2 & \cdots & \mathbf{0}_2 \\ \mathbf{0}_2 & R_{\vec{v}}(\theta_1) & \cdots & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \ddots & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \cdots & R_{\vec{v}}(\theta_{2^{n-1}-1}) \end{pmatrix}$$


**Example:** Outlined block is  $\text{diag}[R_z(\theta_1), R_z(\theta_2), \dots, R_z(\theta_{2^{n-1}})] = \bigwedge_{n-1}^{\text{uni}} [R_z]$  up to SWAP of qubits  $1, n$

**Shende, q-ph/0406176:** **Short** proof of  $2^{n-1}$  CNOTs using induction:  
 $\mathfrak{a}(2^n) = I_2 \otimes \mathfrak{a}(2^{n-1}) \oplus \sigma^z \otimes \mathfrak{a}(2^{n-1})$

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# Universal Circuits

**Goal:** Build a **universal quantum circuit** for  $u$  be  $4^n \times 4^n$  unitary evolution

- Change rotation angles: any  $u$  up to phase
- **Preview:** At least  $4^n - 1$  rotation boxes  $R_{\vec{v}}$ , at least  $\frac{1}{4}(4^n - 3n - 1)$  CNOTs
- Prior art
  - Barenco Bennett Cleve DiVincenzo Margolus Shor Sleator J.Smolin Weinfurter (1995)  $\approx 50n^2 \times 4^n$  CNOTs
  - Vartiainen, Möttönen, Bergholm, Salomaa,  $\approx 8 \times 4^n$  (2003),  $\approx 4^n$  (2004)

## Cosine Sine Decomposition

**Cosine Sine Decomposition:** Any  $v$  a  $2^n \times 2^n$  unitary may be written

$$v = \begin{pmatrix} a_1 & 0 \\ 0 & b_1 \end{pmatrix} \begin{pmatrix} c & -s \\ s & c \end{pmatrix} \begin{pmatrix} a_2 & 0 \\ 0 & b_2 \end{pmatrix} = (a_1 \oplus b_1) \gamma (a_2 \oplus b_2)$$

where  $a_j, b_j$  are  $2^{n-1} \times 2^{n-1}$  unitary,  $c = \sum_{j=0}^{2^{n-1}-1} \cos t_j |j\rangle\langle j|$  and  $s = \sum_{j=0}^{2^{n-1}-1} \sin t_j |j\rangle\langle j|$

- Studied extensively in numerical matrix analysis literature
- **Fast CSD algorithms** exist; reasonable on laptop for  $n = 10$

## Strategy for $\approx 4^n/2$ CNOT Circuit

- Use **CSD** for  $v = (a_1 \oplus b_1)\gamma(c_1 \oplus d_1)$
- Implement  $\gamma = \begin{pmatrix} c & -s \\ s & c \end{pmatrix}$  as **uniformly controlled rotations**
  - uniform control  $\implies$  few CNOTs
- Implement  $a_j \oplus b_j = \begin{pmatrix} a_j & 0 \\ 0 & b_j \end{pmatrix}$  as **quantum multiplexor**
  - Also includes **uniformly controlled rotations**, also inductive
- Induction ends at specialty two-qubit circuit

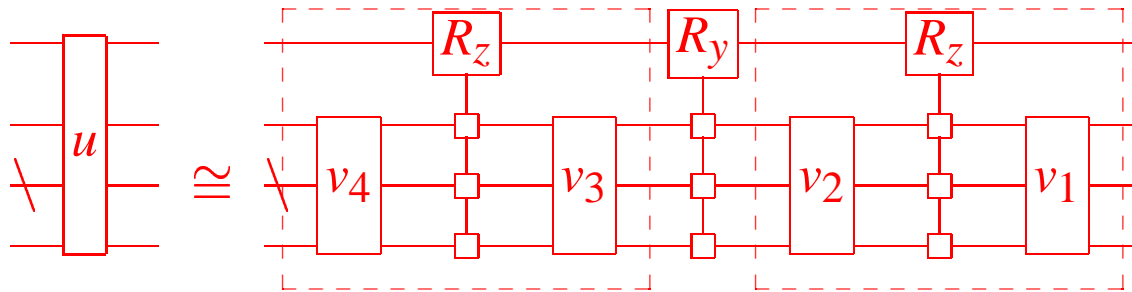
# Quantum Multiplexors

- **Multiplexor:** route computation as control bit 0,1
- $v = a \oplus b$ : Do  $a$  or  $b$  as top qubit  $|0\rangle, |1\rangle$
- **Diagonalization trick:** Solve following system,  $d \in A(2^{j-1})$ ,  
 $u, w$  each some  $2^{n-1} \times 2^{n-1}$  unitary

$$\begin{cases} a &= udw \\ b &= ud^\dagger w \end{cases}$$

- Result:  $a \oplus b = (u \oplus u)(d \oplus d^\dagger)(w \oplus w) = (I_2 \otimes u) \wedge_{n-1}^{\text{uni}}[R_z](I_2 \otimes w)$

## Circuit for $(1/2)$ CNOT per Entry



- Outlined sections are multiplexor implementations
- **Cosine Sine matrix  $\gamma$ :** uniformly controlled  $\wedge_{n-1}^{\text{uni}}[R_y]$ 
  - Only  $2^{n-1}$  CNOTs, converts to  $R_z$  by conjugation by  $HS$

## Circuit Errata

- Lower bound  $\implies$  (can be improved by no more than factor of 2)
- 21 CNOTs in 3 qubits: currently best known
- $\approx 50\%$  CNOTs on bottom two lines
  - Adapts to spin-chain architecture with  $(4.5) \times 4^n$  CNOTs
  - Quantum charge couple device (QCCD) with 3 or 4 qubit chamber?

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# Sard's Theorem

**Def:** A **critical value** of a smooth function of smooth manifolds  $f : M \rightarrow N$  is any  $n \in N$  such that there is some  $p \in M$  with  $f(p) = n$  with the linear map  $(df)_p : T_pM \rightarrow T_nN$  not onto.

**Sard's theorem:** The set of critical values of any smooth map has measure zero.

**Corollary:** If  $\dim M < \dim N$ , then **image(f)** is measure 0.

- $U(2^n) = \{u \in \mathbb{C}^{2^n \times 2^n} ; uu^\dagger = I_{2^n}\}$ : smooth manifold
- Circuit topology  $\tau$  with  $k$  one parameter rotation boxes induces smooth evaluation map  $f_\tau : U(1) \times \mathbb{R}^k \rightarrow U(2^n)$



## Dimension-Based Bounds

- Consequence: Any universal circuit must contain  $4^n - 1$  one parameter rotation boxes
- No consolidation: Boxes separated by at least  $\frac{1}{4}(4^n - 3n - 1)$  CNOTs
  - $\nu$  Bloch sphere rotation:  $\nu = R_x R_z R_x$  or  $\nu = R_z R_x R_z$
  - Diagrams below: consolidation if fewer CNOTs



# On-going Work

- Subgroups  $H$  of unitary group  $U(2^n)$ 
  - More structure, smaller circuits?
  - Symmetries encoded within subgroups  $H$
  - Native gate libraries?
- Special purpose circuits
  - Backwards: quantum circuits for doing numerical linear algebra?
  - Entanglement dynamics and circuit size

<http://www.arxiv.org> **Coordinates**

- Two-qubits: [q-ph/0308045](#)
- Diagonal circuits: [q-ph/0303039](#)
- Uniform control: [q-ph/0404089](#)
- $(1/2)$  CNOT/entry: [q-ph/0406176](#)
- Circuit diagrams by `Qcircuit.tex`: [q-ph/0406003](#)